

Angular Momentum

Introduction

Angular momentum is classically defined as:

$$\mathbf{L} = \mathbf{r} \times \mathbf{p}$$

The quantum-mechanical operator is defined the same way, by just replacing the two quantities with their respective operators.

$$\hat{L}_i = \hat{x}_j \hat{p}_k - \hat{x}_k \hat{p}_j = -i\hbar \left(x_j \frac{\partial}{\partial x_k} - x_k \frac{\partial}{\partial x_j} \right)$$

The three components of angular momentum are not compatible quantities with each other, since they don't commute.

$$[\hat{L}_i, \hat{L}_j] = i\hbar \hat{L}_k \quad \text{Leitrigtharrow} \quad \mathbf{L} \times \mathbf{L} = i\hbar \mathbf{L}$$

In order to fully describe the angular momentum of a quantum-mechanical system we need another quantity, which would be compatible. This quantity is the magnitude of the total angular momentum squared, $\hat{L}^2 = |\hat{\mathbf{L}}|^2$.

Indeed, this quantity commutes with all the components.

$$[\hat{L}^2, \hat{L}_i] = 0$$

Completely describing the angular momentum of a quantum system we use two operators, $\hat{L}^2$ for the magnitude and $\hat{L}_z$ for its projection on the z axis, which is easier to play with in spherical coordinates.

Also, both of those operators commute with the Hamiltonian, and thus have common eigenfunctions. In the end, we have two eigenvalue problems.

$$\hat{L}^2 |l, m\rangle = \hbar^2 l(l+1) |l, m\rangle \quad \hat{L}_z |l, m\rangle = \hbar m |l, m\rangle$$

The eigenvalue $\lambda = l(l+1)$ is written this way because it is easier to handle mathematically, as will be shown later on, without affecting the problem in any way. The quantum numbers l, m are integers and define the so called **Spherical Harmonics** $Y_{lm}(\theta, \phi)$, which are the angular part of the wavefunctions $\psi_{lm}(r, \theta, \phi) = R(r) \cdot Y_{lm}(\theta, \phi)$.

The existence of a central potential $V(\mathbf{r}) = V(-\mathbf{r})$ causes the quantization of the radial part of the wavefunction. Together with spin they form a quartet of quantum numbers (n, l, m_l, m_s) which fully defines the state of a particle around the central potential (e.g. a *single* electron around a nucleus).

Calculating the Eigenvalues

The eigenvalue problem of $\hat{L}_z$ is similar to that of the Hamiltonian for the [Quantum Harmonic Oscillator](#). We are looking for the *raising* and *lowering* operators $\hat{L}_{\pm}$ that will move us from the state $|\ell, m\rangle$ to the state $|\ell, m\pm 1\rangle$.

The only difference is that ℓ should not be affected, meaning that the eigenvalue equation for $\hat{L}^2$ will still hold. ℓ represents the magnitude of the angular momentum, while m its projection on an axis. It is a natural requirement that any action on m , which corresponds to axes rotation, won't affect ℓ in any way.

We start by factoring the magnitude:

$$\hat{L}^2 = \hat{L}_x^2 + \hat{L}_y^2 + \hat{L}_z^2 = \hat{L}_z^2 + (\hat{L}_x + i\hat{L}_y)(\hat{L}_x - i\hat{L}_y)$$

By comparing with the algebra of the harmonic oscillator (factoring $\hat{a}^\dagger \hat{a}$) we practically guess the existence of the operators:

$$\hat{L}_+ = \hat{L}_x + i\hat{L}_y \quad \hat{L}_- = \hat{L}_x - i\hat{L}_y$$

It's obvious that $\hat{L}_{\pm}^\dagger = \hat{L}_{\mp}$. We now need to show that these two operators have the required properties:

- $[\hat{L}_{\pm}, \hat{L}^2] = 0$
- $[\hat{L}_z, \hat{L}_{\pm}] = \pm \hbar \hat{L}_{\pm}$, which leads to $\hat{L}_{\pm} |\ell, m\rangle \propto |\ell, m\pm 1\rangle$

First Property

The $\hat{L}_{\pm}$ operator is a linear combination of $\hat{L}_x, \hat{L}_y$ which commute with $\hat{L}^2$. So $\hat{L}_{\pm}$ commutes too.

$$[\hat{L}_{\pm}, \hat{L}^2] = [\hat{L}_x \pm i\hat{L}_y, \hat{L}^2] = [\hat{L}_x, \hat{L}^2] \pm i[\hat{L}_y, \hat{L}^2] = 0 \pm i \cdot 0 = 0$$

Second Property

$$\begin{aligned} [\hat{L}_z, \hat{L}_{\pm}] &= [\hat{L}_z, \hat{L}_x \pm i\hat{L}_y] = [\hat{L}_z, \hat{L}_x] \pm i[\hat{L}_z, \hat{L}_y] \\ &= \hbar \hat{L}_{\pm} \pm i(-\hbar \hat{L}_{\pm}) = \hbar \hat{L}_{\pm} \mp i\hbar \hat{L}_{\pm} = \pm \hbar \hat{L}_{\pm} \end{aligned}$$

The quantum number m is the eigenvalue of $\hat{L}_z$, so in order to see if $\hat{L}_{\pm}$ indeed takes us from state $|\ell, m\rangle$ to state $|\ell, m\pm 1\rangle$ we can act with it on state $\hat{L}_{\pm} |\ell, m\rangle$.

$$\begin{aligned} \hat{L}_z \hat{L}_{\pm} |\ell, m\rangle &= \hat{L}_{\pm} \hat{L}_z |\ell, m\rangle \pm \hbar \hat{L}_{\pm} |\ell, m\rangle \\ &= \hbar m \hat{L}_{\pm} |\ell, m\rangle \pm \hbar \hat{L}_{\pm} |\ell, m\rangle = \hbar(m \pm 1) \hat{L}_{\pm} |\ell, m\rangle \end{aligned}$$

$$\begin{aligned} \hat{L}_\pm \hbar^{-1} \left| l, m \right\rangle &= \hbar^{-1} \left(m \pm 1 \right) \hat{L}_\pm \left| l, m \right\rangle \\ \end{aligned}$$

The eigenvalue is $\hbar^{-1} \mu = \hbar^{-1} (m \pm 1)$, so the operator $\hat{L}_\pm$ indeed takes us from state $\left| l, m \right\rangle$ to state $\left| l, m \pm 1 \right\rangle$.

Normalization

Accepted Values

Constructing the Eigenfunctions

Solving the Eigenvalue Problem

Using Operators

Example calculations of different quantities

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